Journal de Mathématiques
Mathématiques et Applications
Fondamentales \& Informatique
JMMAFI, ISSN 2411-7188, Vol.6, 2023, pp.21-25


# On Finslerian connections 

Manelo Anona<br>Department of Mathematics and Computer Science Faculty of Science, University of Antananarivo Antananarivo 101, PB: 906, MADAGASCAR e-mail address: mfanona@yahoo.fr


#### Abstract

Given a connection without torsion, we propose a method to know if the connection is Finslerian, in the case of a linear connection if it is Riemannian.


Keywords: Differentiable manifold, Nijenhuis tensor, Riemannian manifold, Spray.
Mathematical Subject Classification (2020): Primary 53XX; Secondary 17B66, 53B05, 53C08.

This article belongs to a special issue on JMMAFI, proceeding of Conférence de Mathématiques en l'honneur de Pr. Emérite Anona M. F. à l'occasion de son 76 ème anniversaire on 15-16-17 March 2023, Université d'Antananarivo, Madagascar, edited by Princy Randriambololondrantomalala.

## 1 Introduction

Let $M$ be a paracompact differentiable manifold of dimension $n \geq 2$ and of class $\mathcal{C}^{\infty}$. The considered connection is an almost product structure [4] which is a version of that of [6] using the formalism of [3].
Let $J$ be the natural tangent structure of tangent bundle $T M, S$ a spray of class $\mathcal{C}^{\infty}$ on $T M-\{0\}, \mathcal{C}^{1}$ on null section, homogeneous of degree $1([C, S]=S), C$ is the Liouville field on the tangent bundle $T M$. The almost product structure $\Gamma=[J, S]$ is a connection without torsion. If $S$ is of class $\mathcal{C}^{2}$ on the null section, then the connection $\Gamma$ is a linear connection without torsion. The curvature of $\Gamma$ is the Nijenhuis tensor of $h, R=\frac{1}{2}[h, h]$, with $h=\frac{I+\Gamma}{2}, I$ being the identity vector 1 -form.

## 2 Finslerian manifold

A map $E$ from $T M-\{0\}$ in $\mathbb{R}^{+}$with $E(0)=0$, of class $\mathcal{C}^{\infty}$ on $\mathcal{T} M$, of class $\mathcal{C}^{1}$ on the null section and homogeneous of degree 2 such that $d d_{J} E$, where $d$ is the exterior derivation,
has a maximal rank, defines a Finslerian structure. The map $E$ is called an energy function. If $E$ is of class $\mathcal{C}^{2}$ on null section, the manifold is Riemannian. The almost product structure, called the canonical connection [4] is given by

$$
\Gamma=[J, S],
$$

where $S$ is a spray defined by [5]

$$
i_{S} d d_{J} E=-d E,
$$

$i_{S}$ being the inner product with respect to $S$. This connection is without torsion and conservative [4].
The fundamental form $\Omega=d d_{J} E$ allows to define, a metric $g$ on the vertical bundle by

$$
g(J X, J Y)=\Omega(J X, Y),
$$

for all $X, Y \in \chi(T M)$, where $\chi(T M)$ denotes the set of all vector fields on $T M$.
There is [4], one and only one metric lift $D$ of a canonical connection such that:

1. $J \mathbb{T}(h X, h Y)=0$,
2. $\mathbb{T}(J X, J Y)=0\left(\mathbb{T}(X, Y)=D_{X} Y-D_{Y} X-[X, Y]\right)$,
3. $D J=0$,
4. $D C=v$,
5. $D \Gamma=0$,
6. $D g=0$.

The linear connection $D$ is called Cartan connection.
Proposition 2.1 ([1]). Let E be an energy function, $\Gamma$ a connection such that $\Gamma=[J, S]$. The following two relationships are equivalent: $i)-i_{S} d d_{J} E=d E$, ii) $d_{h} E=0$.

Remark 2.2 ([4]). This property expresses the fact that the energy function is conserved by parallel transport and the paths of the spray are the solutions of the "Lagrange equations" of energy $E$.

Proposition 2.3. For the connection $\Gamma$ satisfying the proposition 2.1 the scalar $1-$ form $d_{v} E$ is completely integrable.

Proof. The Kernel of $d_{v} E$ is formed by vector fields belonging to the horizontal space Imh $(v \circ h=0)$ and vertical vector fields $J Y$ such that $L_{J Y} E=0, Y \in I m h$, taking into account $v J=J, L_{J Y}$ being the Lie derivative with respect to a vector field $J Y$.
As we have

$$
[h X, h Y]=h[h X, h Y]+v[h X, h Y]=h[h X, h Y]+R(X, Y),
$$

for all $X, Y \in \chi(T M)$, and that $d_{h} E=0$ implies $d_{R} E=0$. We obtain

$$
[h X, h Y] \in \operatorname{Kerd}_{v} E .
$$

It remains to show that $L_{v[h X, J Y]} E=0, \forall X \in \operatorname{Imh}, Y \in \operatorname{Imh}$ satisfying $L_{J Y} E=0$. This is immediate since we have $v=I-h$.

Theorem 2.4. Let $\Gamma=[J, S]$ be a connection. The connection $\Gamma$ comes from a energy function if and only if

1) there is an energy function $E_{0}$ such that $d_{R} E_{0}=0$.
2) the scalar 1-form $d_{v} E_{0}$ is completely integrable.

Then, there exists a constant $\varphi(x)$ on the bundle such that $e^{\varphi(x)} E_{0}$ is the energy function of $\Gamma$.
Proof. This is a consequence of the propositions 2.12 .3 and the proof of theorem 1 of [1].

## 3 Riemannian manifold

In this paragraph, we assure that the energy function $E$ is of class $\mathcal{C}^{2}$ on the null section, then the manifold $(M, E)$ becomes Riemannian. With the Cartan connection $D$, we have [4]

$$
D_{J X} J Y=[J, J Y] X, D_{h X} J Y=[h, J Y] X .
$$

With the linear connection $D$, we associate a curvature

$$
\begin{equation*}
\mathcal{R}(X, Y) Z=D_{h X} D_{h Y} J Z-D_{h Y} D_{h X} J Z-D_{[h X, h Y]} J Z \tag{3.1}
\end{equation*}
$$

for all $X, Y, Z \in \chi(T M)$. The relationship between the curvature $\mathcal{R}$ and $R$ is

$$
\mathcal{R}(X, Y) Z=J[Z, R(X, Y)]-[J Z, R(X, Y)]+R([J Z, X], Y)+R(X,[J Z, Y]) .
$$

for all $X, Y, Z \in \chi(T M)$.
In particular,

$$
\mathcal{R}(X, Y) S=-R(X, Y) .
$$

As the function $E$ is homogeneous of degree 2 and of class $\mathcal{C}^{2}$ null on null section. On an open set $U$ of $M,\left(x^{i}, y^{j}\right) \in T U$, the energy function is written

$$
E=\frac{1}{2} g_{i j}\left(x^{1}, \ldots, x^{n}\right) y^{i} y^{j},
$$

where $g_{i j}\left(x^{1}, \ldots, x^{n}\right)$ are symmetric positive functions such that the matrix $\left(g_{i j}\left(x^{1}, \ldots, x^{n}\right)\right)$ is invertible. And the relation $i_{S} d d_{J} E=-d E$ gives the spray $S$

$$
S=y^{i} \frac{\partial}{\partial x^{i}}-2 G^{i}\left(x^{1}, \ldots, x^{n}, y^{1}, \ldots, y^{n}\right) \frac{\partial}{\partial y^{i}},
$$

we denote

$$
\gamma_{i k j}=\frac{1}{2}\left(\frac{\partial g_{k j}}{\partial x^{i}}+\frac{\partial g_{i k}}{\partial x^{j}}-\frac{\partial g_{i j}}{\partial x^{k}}\right)
$$

and

$$
\gamma_{i j}^{k}=g^{k l} \gamma_{i l j},
$$

we have

$$
G^{k}=\frac{1}{2} y^{i} y^{j} \gamma_{i j}^{k} .
$$

We note $\Gamma_{i}^{j}(x, y)=y^{l} \Gamma_{i l}^{j}(x)$, the horizontal projector is written

$$
\left\{\begin{array}{l}
h\left(\frac{\partial}{\partial x^{i}}\right)=\frac{\partial}{\partial x^{i}}-\Gamma_{i}^{j} \frac{\partial}{\partial y^{j}} \\
h\left(\frac{\partial}{\partial y^{j}}\right)=0
\end{array}\right.
$$

The vertical projector becomes

$$
\left\{\begin{array}{l}
v\left(\frac{\partial}{\partial x^{i}}\right)=\Gamma_{i}^{j} \frac{\partial}{\partial y^{j}} \\
v\left(\frac{\partial}{\partial y^{j}}\right)=\frac{\partial}{\partial y^{j}}
\end{array} .\right.
$$

The curvature $R=\frac{1}{2}[h, h]$ is then

$$
R=\frac{1}{2} R_{i j}^{k} d x^{i} \wedge d x^{j} \otimes \frac{\partial}{\partial y^{k}} \text { with } R_{i j}^{k}=\frac{\partial \Gamma_{i}^{k}}{\partial x^{j}}-\frac{\partial \Gamma_{j}^{k}}{\partial x^{i}}+\Gamma_{i}^{l} \frac{\partial \Gamma_{j}^{k}}{\partial y^{l}}-\Gamma_{j}^{l} \frac{\partial \Gamma_{i}^{k}}{\partial y^{l}}, i, j, k, l \in\{1, \ldots, n\} .
$$

Proposition 3.1 ([2]). On a Riemannian manifold ( $M, E$ ), the horizontal nullity space of the curvature $R$ is generated as a module by the projectable vector fields belonging to this nullity space and, orthogonal to the image space $\operatorname{ImR}$ of the curvature $R$ and $h N_{R}=h N_{\mathcal{R}}$.

Proposition 3.2. On a Riemannian manifold ( $M, E$ ), the horizontal space and the space $\operatorname{Im} R$ of the curvature $R$ generate a Lie algebra on $\mathcal{F}(T M)$ whose the vertical space is orthogonal to $J N_{R}+\langle C\rangle_{\mathcal{F}(T M)}$.

Proof. It is immediate to see that $\operatorname{Im} R$ is orthogonal to $J N_{R}+\langle C\rangle_{\mathcal{F}(T M)}$, from the relations between the curvature $R, \mathcal{R}$ and $h N_{R}$ [1].
Let be $X, Y \in \chi(T M)$, we have

$$
[h X, h Y]=h[h X, h Y]+v[h X, h Y]=h[h X, h Y]+R(X, Y),
$$

ie, $[h X, h Y] \in I m h+I m R$, for all $X, Y \in \chi(T M)$.
Let be $X \in h N_{R}$ and $J Y \perp J N_{R}+\{C\}$. We have $g(J Y, J X)=0$. We can consider $X$ projectable according to the proposition 3.1 since $g$ is bilinear. Given $D g=0$, we can write

$$
D_{h Z} g(J Y, J X)-g\left(D_{h Z} J Y, J X\right)-g\left(J Y, D_{J Z} J X\right)=0 .
$$

From the relation $h^{2}=h$, we get

$$
D_{h Z} J Y=v[h Z, J Y] \text { and } D_{h Z} J X=v[h Z, J X]=[h, J X] h Z=0 \text {, }
$$

according to the proposition 4 [1]. The relation $v[h Z, J Y]$ is orthogonal to $J N_{R}$. Similarly, $v[h Z, J Y]$ is orthogonal to $C=J S$, given $[C, h]=0$. If $J Y, J Z$ are orthogonal to $J N_{R}$, we have $g(J Y, J X)=0, g(J Z, J X)=0$ for all $X \in h N_{R}$.
From the relation $D g=0$, we can write $g\left(D_{h Z} J Y, J X\right)+g\left(J Y, D_{J Z} J X\right)=0$ and $D_{J Z} J Y=$ $[J, J Y] Z=J[J Z, Y]$, taking into account $[J, J]=0$. Likewise, we find $D_{J Z} J X=J[J Z, X]$. We then have $g([J Z, J Y], J X)=g(J[J Z, Y]+J[Z, J Y], J X)=0$. That is, $[J Z, J Y]$ is orthogonal to $J N_{R}$. If $J Y, J Z$ are orthogonal to $C$, it means $L_{J Y} E=0$ and $L_{J Z} E=0$. We then have $L_{[J Y, J Z]} E=0$. Hence the result.

Remark 3.3. For a Riemannian manifold, condition 2) of the theorem 2.4 could become $H+\operatorname{ImR}$ completely integrable by the proposition 3.2 if $h N_{R}=\{0\}$.
An energy function is written

$$
E_{0}=\frac{1}{2} g_{i j}^{0} y^{i} y^{j}
$$

Thus, the relation $d_{R} E=0$ is equivalent to the following system of equations

$$
\left\{\begin{array}{l}
g_{k l}^{0} R_{l, i j}^{k}=0 \\
g_{k l}^{0} R_{r, i j}^{k}=-g_{k r}^{0} R_{l, i j}^{k} \quad \text { with } l \neq r .
\end{array}\right.
$$

In the matrix form, the system is written

$$
\left(\begin{array}{ccc}
g_{11}^{0} & \ldots & g_{n 1}^{0} \\
\vdots & \ddots & \vdots \\
g_{1 n}^{0} & \ldots & g_{n n}^{0}
\end{array}\right)\left(\begin{array}{ccc}
R_{1, i j}^{1} & \ldots & R_{n, i j}^{1} \\
\vdots & \ddots & \vdots \\
R_{1, i j}^{n} & \ldots & R_{n, i j}^{n}
\end{array}\right)=\left(\begin{array}{cccc}
0 & & & \\
\ddots & -{ }^{t} \mathbf{A} \\
\mathbf{A} & \ddots & \\
& & & 0
\end{array}\right) .
$$

Remark 3.4. In [7], an answer was given. Let $M$ be a paracompact differentiable manifold, and $D$ a linear connection on $M$, without torsion. For $D$ to come from a Riemannian structure, it is necessary and sufficient that its groups of holonomy are relatively compact.

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